Electron Interactions and Nanotube Fluorescence Spectroscopy

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Large radius theory of optical transitions in semiconducting nanotubes derived from low energy theory of graphene Phys. Rev. Lett. in press cond-mat/ 0403153

- Brief Introduction to nanotubes
- Independent electron model for optical spectra
- 2D interactions: nonlinear scaling with 1/R
- 1D interactions: excitons
- Short Range Interactions: exciton fine structure





Carbon Nanotubes as Electronic Materials



A Molecular Quantum Wire

Tans et al. (Nature 1998)

- Ballistic Conductor
- Field Effect Transistor
- Logic Gates

Carbon Nanotubes as Optical Materials

Photoluminescence



- Nanotubes in surfactant micelles Bachillo et al. (2002).
- Photoluminescence from individual suspended nanotubes Lefebvre et al. (2003).

Electroluminescence & Photoconductivity



• Infrared Emission, photoconductivity in individual nanotube field effect devices Freitag et al., (IBM) 2004

Carbon Nanotube : Wrapped Graphene





 $\mathbf{C} = \mathbf{n}_1 \ \mathbf{a}_1 + \mathbf{n}_2 \ \mathbf{a}_2$



Tubes characterized by $[n_1,n_2]$ or

- Radius : $\mathbf{R} = |\mathbf{C}|/2\pi$
- Chiral Angle : $0 < \theta < 30^{\circ}$
- Chiral Index : $v = n_1 n_2 \mod 3$ = 0,1,-1

Electronic Structure

Metal

• Finite Density of States (DOS) at Fermi Energy

Semiconductor

• Gap at Fermi Energy

Graphene

- Zero Gap Semiconductor
- Zero DOS metal





Tight Binding model: $\hbar v_F = \sqrt{3}\gamma_0 a / 2$; $\gamma_0 = 2.5 \text{ eV}$

Wrap it up.....

• Flat Graphene: A zero gap semiconductor



• Periodic boundary conditions on cylinder:



 $n_1 - n_2 = +/-1 \mod 3$ Semiconductor



Near-infrared Photoluminescence from Single-wall Carbon Nanotubes

O'Connel et al. (Science 02) Bachillo et al. (Science 02)



Excitation (661 nm)

Emission (> 850 nm)

Nanotube Fluorescence Spectroscopy ^{O'Connel et al. (Science 02)} Bachillo et al. (Science 02)



Each peak in the correlation plot corresponds to a particular species $[n_1,n_2]$ of semiconducting nanotube

GOAL:

Understand observed transition energies in terms of low energy properties of an ideal 2 dimensional graphene sheet.

Free Electron Theory of Nanotube Bandgaps

Systematic expansion for large radius, R

• Zeroth order:

$$E_n^0 = \frac{2\hbar v_F}{3} \frac{n}{R} \quad (n = 1, 2, 4, 5, ...)$$



Free Electron Theory of Nanotube Bandgaps

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• Trigonal Warping Correction

$$\Delta E_n^{T.W.} \propto (-1)^n v \frac{n^2 \sin 3\theta}{R^2}$$



• Curvature Correction $\Delta E_n^C \propto (-1)^n \nu \frac{\sin 3\theta}{R^2}$





Curvature and Trigonal Warping:

- Vary as $1/R^2$
- Alternate with band index n
- Alternate with chiral index v
- Vanish for armchair tubes, $\theta=0$

Different dependence on n

The large R limit is most accurate for nearly armchair tubes: $\theta \sim 0$ E_n⁰ describes tight binding gaps accurately for R > .5 nm

Nanotube Assignments from Pattern of sin $3\theta/R^2$ Deviations

Experimental "Ratio Plot"

deviations) С D 2.2 2.0 V22/ V11 v_{22}/v_{11} 1.8 2.0 1.6 1.6 1.4 1.2 900 500 600 800 1000 1100 500 600 700 800 700 900 1200 Excitation wavelength (nm) Excitation wavelength (nm)

- By comparing the experimental and theoretical ratio plots the [n₁,n₂] values (and hence R and θ) for each peak can be identified.
- Corroborated by Raman spectroscopy of the radial breathing mode.

Theory (includes $\sin 3\theta/R^2$

The Ratio Problem

- Free electron theory predicts $E_{22} / E_{11} \rightarrow 2$ for $R \rightarrow \infty$
- Consequence of linear dispersion of graphene



Scaling of Optical Transition Energies

(Kane,Mele '04)



- Free electrons for $\theta = 0$ $E_{nn}^0(R) = 2\hbar v_F n / 3R$
- $v \sin 3\theta / R^2$ deviations are clear
- Separatrix between v=+1 and v=-1describes nearly armchair tubes with θ=0, where sin 3θ/R² deviations vanish.

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Scaling of Optical Transition Energies

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- Free electrons for $\theta=0$ $E_{nn}^{0}(R) = 2\hbar v_{F}n/3R$ Ratio Problem: $E_{22}/E_{11} < 2$ Blue Shift Problem: $E_{nn}(R) > E_{nn}^{0}(R)$ Worse for large R

Nonlinear scaling $E_{nn}(R)=E(q_n=n/3R)$ accounts for both effects.

Electron Interactions in large radius tubes

For $2\pi R$ >>a electron interactions can be classified into three regimes, which lead to distinct physical effects.

• Long Range Interaction : $(r > 2\pi R)$

One Dimensional in character

- Strongly bound excitons
- Intermediate Range Interaction : $(a < r < 2\pi R)$

Two Dimensional in character

- Nonlinear Scaling with n/R
- Short Range Interaction : (r~a)

Atomic in character









Long Range Interaction : (r > $2\pi R$) $V(z) = \frac{e^2}{\varepsilon |z|}$

- Renormalize Single Particle Gap Increase observed energy gap
- Leads to exciton binding
 Decrease observed energy gap



- Gap renormalization and exciton binding largely cancel each other.

Cancellation of gap renormalization and exciton binding:

 Single Particle excitation: Self energy ~ e²/εR Depends on dielectric environment



 Particle-hole excitation: Bound exciton is unaffected by the long range part of the interaction.



The cancellation is exact for an infinite range interaction

• Coulomb Blockade Model :

Bare gap: 2Δ Interaction energy: U N²/2

- Single particle gap $2\Delta + U$
- Particle-hole gap 2Δ

Intermediate Range Interaction: (a < r < $2\pi R$)

- Leads to nonlinear q log q dispersion of graphene.
- Responsible for nonlinear scaling of $E_{11}(n/R)$.



Short Range Interaction: (r ~ a)

- Leads to "fine structure" in the exciton spectrum: S=0,1, etc.
- Splittings ~ e² a / R²

Interactions in 2D Graphene Gonzalez, Guinea, Vozmediano, PRB 99

$$H = \hbar \mathbf{v}_{\mathbf{F}} \int d^2 r \psi^{\dagger} \frac{\boldsymbol{\sigma} \cdot \nabla}{i} \psi + \boldsymbol{e}^2 \int d^2 r d^2 r' \frac{n(r)n(r')}{2|r-r'|}$$

• Renormalized Quasiparticle Dispersion:

Singularity due to long range Coulomb interaction $V(q) = 2\pi e^2/q$.

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• "Dielectric Screening" in 2 Dimensions

 $g_{screened} = g/\epsilon$ $\Pi_{static}(q) = q/4v_F$ $\epsilon_{static} = 1+g\pi/2$

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• Renormalized Quasiparticle Dispersion:

$$E(q) = \hbar v_F q \left(1 + \frac{g}{4} \log \frac{\Lambda}{q} \right) \qquad g = e^2 / \hbar v_F$$

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• "Dielectric Screening" in 2 Dimensions

 $g_{\text{screened}} = g/\epsilon$ $\Pi_{\text{static}}(q) = q/4v_F$ $\epsilon_{\text{static}} = 1+g\pi/2$

• Scaling Theory $g = g(\Lambda)$; $v_F = v_F(\Lambda)$

$$\frac{dg}{d\ln\Lambda} = -\frac{1}{4}g^2$$

 $\frac{1}{4}g^2$ Marginally Irrelevant

Marginal Fermi Liquid

• q ln q correction is exact for q#0

Compare 2D Theory with Experiment



Free electron Theory2D Interacting Theory

$$E_{nn} = E(q_n = n/3R)$$
$$E(q) = 2\hbar v_F q \left(1 + \frac{g}{4} \log \frac{\Lambda}{q}\right)$$

$$\hbar v_F = .47 \text{ eV nm}$$

 $g = \frac{e^2}{\epsilon \hbar v_F} = 1.2 \implies \epsilon \sim 2.5$

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The optical spectra reflects the finite size scaling of the 2D Marginal Fermi Liquid

Exciton effects: Compute particle-hole binding due to statically screened interaction (similar to Ando '97).



- Lowest exciton dominates oscillator strength for each subband.
- Lineshape for absorption is not that of van Hove singularity.
- Large bandgap renormalization mostly cancelled by exciton binding.

Scaling behavior: $E_n(R) = E(q_n \stackrel{?}{=} n/3R)$



Exciton Fine Structure



Degenerate exciton states:
e :
$$k = K$$
 or K' ; $s = \uparrow$ or \downarrow 16

h:
$$k = K \text{ or } K'$$
; $s = \uparrow \text{ or } \downarrow \int \text{ states}$

Degeneracy lifted by short range $(q \sim 1/a)$ interactions:



Exciton Eigenstates:

Classify by momentum, spin, parity under C₂ rotation



See also Zhao, Mazumdar PRL 04

Conclusion

Fluorescence spectroscopy data for nearly armchair tubes is well described by a systematic large radius theory.

- 2D interactions:
 - q log q renormalization of graphene dispersion.
 - Non linear scaling with 1/R.
 - Explains ratio problem and blue shift problem.
- 1D interactions
 - Lead to large gap enhancement AND large exciton binding
 - Largely cancels in optical experiments revealing 2D effects.
- Short Range interactions
 - -Lead to fine structure in exciton levels
 - -Dark Ground State
- Experiments: measure single particle energy gap
 - Tunneling (complicated by screening)
 - Photoconductivity
 - Activated transport